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Weight bases of Gelfand–Tsetlin type for representations of classical Lie algebras

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Abstract. This paper completes a series devoted to explicit constructions of finite-dimensional irreducible representations of the classical Lie algebras. Here the case of odd orthogonal Lie algebras (of type B) is considered (two previous papers dealt with C and D types). A weight basis for each representation of the Lie algebra $\mathfrak{o}(2n+1)$ is constructed. The basis vectors are parametrized by Gelfand–Tsetlin-type patterns. Explicit formulae for the matrix elements of generators of $\mathfrak{o}(2n+1)$ in this basis are given. The construction is based on the representation theory of the Yangians.

1. Introduction

In this paper we give an explicit construction of each finite-dimensional irreducible representation V of an odd orthogonal Lie algebra $\mathfrak{o}(2n+1)$ (i.e. a simple complex Lie algebra of type B). A *weight* basis in V is obtained by the application of certain elements of the enveloping algebra (the lowering operators) to the highest weight vector. Explicit formulae for the matrix elements of generators of the Lie algebra $\mathfrak{o}(2n+1)$ in this basis are given. We follow an approach applied in the previous papers [11] and [12] where similar results were obtained for the C and D type Lie algebras.

Let $\mathfrak{g}_n$ denote the rank n simple Lie algebra of type A, B, C, or D. The restriction of a finite-dimensional irreducible representation V of $\mathfrak{g}_n$ to the subalgebra $\mathfrak{g}_{n-1}$ is multiplicity-free for the A type case (see below), and it is not necessarily so for the B, C, D types. Gelfand and Tsetlin [5] used the chain of subalgebras

$$\mathfrak{g}_1 \subset \mathfrak{g}_2 \subset \cdots \subset \mathfrak{g}_n \tag{1.1}$$

to parametrize basis vectors in V and give formulae for the matrix elements of generators for the A type case. Different approaches to derive these formulae are used, e.g., in [7,9,14,15,21,22]. More references and a discussion of the history of the problem of bases for representations of the classical Lie algebras can be found in [11].

Analogous results for representations of the orthogonal Lie algebra $\mathfrak{o}(N)$ are obtained by Gelfand and Tsetlin in [6]. Derivations of the matrix element formulae are given in [8,18,20]. Here the chain (1.1) is replaced with one which involves the orthogonal Lie algebras of both series B and D:

$$\mathfrak{o}(2) \subset \mathfrak{o}(3) \subset \cdots \subset \mathfrak{o}(N). \tag{1.2}$$

However, these inclusions are not regular, i.e., they are not consistent with the corresponding root systems. As a consequence of that, the basis vectors lose the weight property which holds in the A type case. They are not eigenvectors for the elements of the Cartan subalgebra.

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To get a weight basis we propose to use the chain (1.1) for the B, C, D types as well. We 'separate' the multiplicities occurring in the reduction $\mathfrak{g}_n \downarrow \mathfrak{g}_{n-1}$ by applying the representation theory of the Yangians. Namely, the subspace V_μ^+ of $\mathfrak{g}_{n-1}$-highest vectors of weight μ in V possesses a natural structure of a representation of the *twisted Yangian* $\mathrm{Y}^+(2)$ or $\mathrm{Y}^-(2)$, in the orthogonal and symplectic cases, respectively. The twisted Yangians are introduced and studied by Olshanski [17]; see also [13] for a detailed exposition. The action of $\mathrm{Y}^{\pm}(2)$ in the space V_μ^+ arises from his *centralizer construction* [17]. Finite-dimensional irreducible representations of the twisted Yangians are classified in [10]. In particular, it turns out that the representation V_μ^+ of $\mathrm{Y}^{\pm}(2)$ can be extended to a larger algebra, the *Yangian* $\mathrm{Y}(2)$ for the Lie algebra $\mathfrak{gl}(2)$. The algebra $\mathrm{Y}(2)$ and its representations are very well studied; see [2, 19]. In particular, a large class of representation of $\mathrm{Y}(2)$ admits Gelfand–Tsetlin-type bases associated with the inclusion $\mathrm{Y}(1) \subset \mathrm{Y}(2)$; see [9, 16]. This allows us to get a natural basis in the space V_μ^+, and then by induction to get a basis in the entire space V.

Note that in the case of C or D type the $\mathrm{Y}(2)$-module V_μ^+ is irreducible while in the B type case it is a direct sum of two irreducible submodules. This does not lead, however, to major differences in the constructions, and the final formulae are similar in all the three cases.

Our calculations of the matrix elements of the generators of $\mathfrak{g}_n$ are based on the relationship between the twisted Yangian $\mathrm{Y}^{\pm}(2)$ and the *transvector algebra* $\mathrm{Z}(\mathfrak{g}_n, \mathfrak{g}_{n-1})$ (it is also called the *Mickelsson algebra* or *S-algebra*). It is generated by the *raising* and *lowering operators* which preserve the subspace V^+ of $\mathfrak{g}_{n-1}$-highest vectors in V. The algebraic structure of the transvector algebras is studied in detail in [23] with the use of the extremal projections for reductive Lie algebras [1]. We construct an algebra homomorphism $\mathrm{Y}^{\pm}(2) \to \mathrm{Z}(\mathfrak{g}_n, \mathfrak{g}_{n-1})$ which allows us to express the generators of the twisted Yangian, as operators in V_μ^+, in terms of the raising and lowering operators.

Explicit combinatorial constructions of the fundamental representations of the symplectic and odd orthogonal Lie algebras were recently given by Donnelly [3]. He also showed that in the symplectic case the basis of [11] for the fundamental representations coincides, up to a scaling, with a basis of his [4]. It is likely that a similar connection exists in the odd orthogonal case.

We conclude the introduction with a brief discussion of the A type case. We outline a construction of the Gelfand–Tsetlin basis which is based on the well known technique of the lowering and raising operators; see, e.g., [22]. We shall employ this technique in the B case as well. During the discussion we introduce some notation to be used in the following.

Let E_{ij}, $i, j = 1, \dots, n$ denote the standard basis of the general linear Lie algebra $\mathfrak{gl}(n)$ over the field of complex numbers. The subalgebra $\mathfrak{gl}(n-1)$ is spanned by the basis elements E_{ij} with $i, j = 1, \dots, n-1$. We denote by $L(\lambda)$ the finite-dimensional irreducible representation of $\mathfrak{gl}(n)$ with the highest weight $\lambda = (\lambda_1, \dots, \lambda_n)$. Here the λ_i are complex parameters such that all the differences $\lambda_i - \lambda_{i+1}$ are non-negative integers. The representation $L(\lambda)$ contains a unique, up to a multiple, nonzero vector ξ (the highest vector) such that $E_{ii}\xi = \lambda_i \xi$ for all i and $E_{ij}\xi = 0$ for $1 \leqslant i < j \leqslant n$. Given a $\mathfrak{gl}(n-1)$-highest weight $\mu = (\mu_1, \dots, \mu_{n-1})$ we denote by $L(\lambda)_\mu^+$ the subspace in $L(\lambda)$ of $\mathfrak{gl}(n-1)$-highest vectors of weight μ. It is well known [21] that the space $L(\lambda)_\mu^+$ is either trivial or one-dimensional. Moreover, $\dim L(\lambda)_\mu^+ = 1$ if and only if

$$\lambda_i - \mu_i \in \mathbb{Z}_+ \qquad \text{and} \qquad \mu_i - \lambda_{i+1} \in \mathbb{Z}_+ \qquad \text{for} \quad i = 1, \dots, n-1. \tag{1.3}$$

In other words, the restriction of $L(\lambda)$ to the subalgebra $\mathfrak{gl}(n-1)$ is multiplicity-free:

$$L(\lambda)|_{\mathfrak{gl}(n-1)} \simeq \bigoplus L'(\mu) \tag{1.4}$$

where $L'(\mu)$ is the irreducible $\mathfrak{gl}(n-1)$-module with the highest weight μ satisfying the conditions (1.3). Further restrictions of $L(\lambda)$ to the subalgebras of the chain (1.1) yield a decomposition of $L(\lambda)$ into a direct sum of one-dimensional subspaces parametrized by the *Gelfand–Tsetlin patterns* Λ. These are arrays of row vectors

$$\begin{array}{ccccccc}
\lambda_{n1} & & \lambda_{n2} & & \cdots & & \lambda_{nn} \\
& \lambda_{n-1,1} & & \cdots & & \lambda_{n-1,n-1} & \\
& \cdots & & \cdots & & \cdots & \\
& & \lambda_{21} & & \lambda_{22} & & \\
& & & \lambda_{11} & & &
\end{array}$$

such that the upper row coincides with λ and the following conditions hold:

$$\lambda_{ki} - \lambda_{k-1,i} \in \mathbb{Z}_+ \qquad \lambda_{k-1,i} - \lambda_{k,i+1} \in \mathbb{Z}_+ \qquad i = 1, \ldots, k-1 \tag{1.5}$$

for each $k = 2, \ldots, n$. Corresponding basis vectors are constructed by means of the *lowering operators* originated from [14]. Set for $i = 1, \ldots, k-1$

$$z_{ki} = \sum_{i<i_1<\cdots<i_s<k} E_{i_1 i} E_{i_2 i_1} \cdots E_{i_s i_{s-1}} E_{k i_s} (h_i - h_{j_1}) \cdots (h_i - h_{j_r}) \tag{1.6}$$

where $h_i = E_{ii} - i + 1$ and $\{j_1, \ldots, j_r\}$ is the complement to the subset $\{i_1, \ldots, i_s\}$ in $\{i+1, \ldots, k-1\}$. Set

$$\xi_\Lambda = \prod_{k=2,\ldots,n}^{\rightarrow} (z_{k1}^{\lambda_{k1}-\lambda_{k-1,1}} \cdots z_{k,k-1}^{\lambda_{k,k-1}-\lambda_{k-1,k-1}})\xi.$$

The vectors $\{\xi_\Lambda\}$ form a basis of $L(\lambda)$ and the action of generators of $\mathfrak{g}_n$ is given by the formulae

$$E_{kk}\xi_\Lambda = \left(\sum_{i=1}^{k} \lambda_{ki} - \sum_{i=1}^{k-1} \lambda_{k-1,i} \right) \xi_\Lambda$$

$$E_{k,k+1}\xi_\Lambda = -\sum_{i=1}^{k} \frac{(l_{ki} - l_{k+1,1}) \cdots (l_{ki} - l_{k+1,k+1})}{(l_{ki} - l_{k1}) \cdots \wedge \cdots (l_{ki} - l_{kk})} \xi_{\Lambda+\delta_{ki}}$$

$$E_{k+1,k}\xi_\Lambda = \sum_{i=1}^{k} \frac{(l_{ki} - l_{k-1,1}) \cdots (l_{ki} - l_{k-1,k-1})}{(l_{ki} - l_{k1}) \cdots \wedge \cdots (l_{ki} - l_{kk})} \xi_{\Lambda-\delta_{ki}}.$$

Here $l_{ki} = \lambda_{ki} - i + 1$ and the arrays $\Lambda \pm \delta_{ki}$ are obtained from Λ by replacing λ_{ki} by $\lambda_{ki} \pm 1$. It is supposed that $\xi_\Lambda = 0$ if the array Λ is not a pattern; the symbol $\wedge$ indicates that the zero factor in the denominator is skipped.

The matrix element formulae are derived with the use of the *raising operators* defined by analogy with (1.6); see, e.g., [22] for details. Some other derivations of the formulae are given, e.g., in [7, 9, 15]. Note that the original basis [5] is orthonormal. The basis vectors there coincide with the ξ_Λ, up to a norm factor, which can be explicitly calculated; cf [22].

2. A basis for odd orthogonal Lie algebras

We shall enumerate the rows and columns of $(2n+1) \times (2n+1)$ matrices over $\mathbb{C}$ by the indices $-n, \ldots, -1, 0, 1, \ldots, n$.

2.1. Main theorem

We keep the notation $E_{ij}, i, j = -n, \ldots, n$ for the standard basis of the Lie algebra $\mathfrak{gl}(2n+1)$. Introduce the elements

$$F_{ij} = E_{ij} - E_{-j,-i}. \tag{2.1}$$

We have $F_{-j,-i} = -F_{ij}$. In particular, $F_{-i,i} = 0$ for all i. The orthogonal Lie algebra $\mathfrak{g}_n := \mathfrak{o}(2n+1)$ can be identified with the subalgebra in $\mathfrak{gl}(2n+1)$ spanned by the elements F_{ij}, $i, j = -n, \ldots, n$. The subalgebra $\mathfrak{g}_{n-1}$ is spanned by the elements (2.1) with the indices i, j running over the set $\{-n+1, \ldots, n-1\}$. Denote by $\mathfrak{h} = \mathfrak{h}_n$ the diagonal Cartan subalgebra in $\mathfrak{g}_n$. The elements $F_{11}, \ldots, F_{nn}$ form a basis of $\mathfrak{h}$.

The finite-dimensional irreducible representations of $\mathfrak{g}_n$ are in a one-to-one correspondence with n-tuples $\lambda = (\lambda_1, \ldots, \lambda_n)$ where all the entries λ_i are simultaneously integers or half-integers (elements of the set $\frac{1}{2} + \mathbb{Z}$) and the following inequalities hold:

$$0 \geqslant \lambda_1 \geqslant \lambda_2 \geqslant \cdots \geqslant \lambda_n.$$

Such an n-tuple λ is called the highest weight of the corresponding representation which we shall denote by $V(\lambda)$. It contains a unique, up to a multiple, nonzero vector ξ (the highest vector) such that $F_{ii}\xi = \lambda_i \xi$ for $i = 1, \ldots, n$ and $F_{ij}\xi = 0$ for $-n \leqslant i < j \leqslant n$. Denote by $V(\lambda)^+$ the subspace of $\mathfrak{g}_{n-1}$-highest vectors in $V(\lambda)$:

$$V(\lambda)^+ = \{\eta \in V(\lambda) | F_{ij}\eta = 0 \quad -n < i < j < n\}.$$

Given a $\mathfrak{g}_{n-1}$-highest weight $\mu = (\mu_1, \ldots, \mu_{n-1})$ we denote by $V(\lambda)^+_\mu$ the corresponding weight subspace in $V(\lambda)^+$:

$$V(\lambda)^+_\mu = \{\eta \in V(\lambda)^+ | F_{ii}\eta = \mu_i \eta \quad i = 1, \ldots, n-1\}.$$

By the branching rule for the reduction $\mathfrak{g}_n \downarrow \mathfrak{g}_{n-1}$ [21] we have

$$V(\lambda)|_{\mathfrak{g}_{n-1}} \simeq \bigoplus c(\mu) V'(\mu) \tag{2.2}$$

where $V'(\mu)$ is the irreducible finite-dimensional representation of $\mathfrak{g}_{n-1}$ with the highest weight μ, and $c(\mu)$ equals the number of n-tuples $\rho = (\rho_1, \ldots, \rho_n)$ satisfying the inequalities

$$\begin{aligned} -\lambda_1 \geqslant \rho_1 \geqslant \lambda_1 \geqslant \rho_2 \geqslant \lambda_2 \geqslant \cdots \geqslant \rho_{n-1} \geqslant \lambda_{n-1} \geqslant \rho_n \geqslant \lambda_n \\ -\mu_1 \geqslant \rho_1 \geqslant \mu_1 \geqslant \rho_2 \geqslant \mu_2 \geqslant \cdots \geqslant \rho_{n-1} \geqslant \mu_{n-1} \geqslant \rho_n \end{aligned} \tag{2.3}$$

with all the ρ_i and μ_i being simultaneously integers or half-integers together with the λ_i. Any nonzero vector in $V(\lambda)^+_\mu$ generates a $\mathfrak{g}_{n-1}$-submodule in $V(\lambda)$ isomorphic to $V'(\mu)$. We obviously have $\dim V(\lambda)^+_\mu = c(\mu)$. Basis vectors in $V(\lambda)^+_\mu$ can be parametrized by the n-tuples ρ. We shall be using an equivalent parametrization by $(n+1)$-tuples $\nu = (\sigma, \nu_1, \ldots, \nu_n)$, where $\nu_i = \rho_i$ for $i \geqslant 2$, and

$$(\sigma, \nu_1) = \begin{cases} (0, \rho_1) & \text{if} \quad \rho_1 \leqslant 0 \\ (1, -\rho_1) & \text{if} \quad \rho_1 > 0. \end{cases}$$

A parametrization of basis vectors in $V(\lambda)$ is obtained by using its subsequent restrictions to the subalgebras of the chain $\mathfrak{g}_1 \subset \mathfrak{g}_2 \subset \cdots \subset \mathfrak{g}_{n-1} \subset \mathfrak{g}_n$. Define a *pattern* Λ associated with λ as an array of the form

$$\begin{array}{llllllllll}
\sigma_n & & \lambda_{n1} & & \lambda_{n2} & & \cdots & & \lambda_{nn} \\
& \lambda'_{n1} & & \lambda'_{n2} & & \cdots & & \lambda'_{nn} & \\
\sigma_{n-1} & & \lambda_{n-1,1} & & \cdots & & \lambda_{n-1,n-1} & & \\
& \lambda'_{n-1,1} & & \cdots & & \lambda'_{n-1,n-1} & & & \\
& \cdots & & \cdots & & \cdots & & & \\
\sigma_1 & & \lambda_{11} & & & & & & \\
& \lambda'_{11} & & & & & & &
\end{array}$$

such that $\lambda = (\lambda_{n1}, \ldots, \lambda_{nn})$, each σ_k is 0 or 1, the remaining entries are all nonpositive integers or nonpositive half-integers together with the λ_i, and the following inequalities hold:

$$\lambda'_{k1} \geqslant \lambda_{k1} \geqslant \lambda'_{k2} \geqslant \lambda_{k2} \geqslant \cdots \geqslant \lambda'_{k,k-1} \geqslant \lambda_{k,k-1} \geqslant \lambda'_{kk} \geqslant \lambda_{kk}$$

for $k = 1, \ldots, n$; and

$$\lambda'_{k1} \geqslant \lambda_{k-1,1} \geqslant \lambda'_{k2} \geqslant \lambda_{k-1,2} \geqslant \cdots \geqslant \lambda'_{k,k-1} \geqslant \lambda_{k-1,k-1} \geqslant \lambda'_{kk}$$

for $k = 2, \ldots, n$; in addition, in the case of integer λ_i the condition

$$\lambda'_{k1} \leqslant -1 \qquad \text{if} \quad \sigma_k = 1$$

should hold for all $k = 1, \ldots, n$. Let us set $l_{k0} = -\frac{1}{2}$ for all k, and

$$l_{ki} = \lambda_{ki} - i + \tfrac{1}{2} \qquad l'_{ki} = \lambda'_{ki} - i + \tfrac{1}{2} \qquad 1 \leqslant i \leqslant k \leqslant n.$$

Given a pattern Λ set for $i = 0, 1, \ldots, k-1$

$$A_{ki} = \prod_{a=1, a\neq i}^{k-1} \frac{1}{l_{k-1,i} - l_{k-1,a}} \cdot \prod_{a=1}^{k-1} \frac{1}{l_{k-1,i} + l_{k-1,a}}.$$

Furthermore, introduce polynomials $B_{ki}(x)$ by

$$B_{ki}(x) = \prod_{a=1, a\neq i}^{k} \frac{(x + l'_{ka} + 1)(x - l'_{ka})}{l'_{ka} - l'_{ki}}$$

and define the numbers C_{ki} by

$$C_{ki} = l'_{ki}(1 - 2\sigma_k - 2l'_{ki}) \prod_{a=1}^{k} (l_{ka} - l'_{ki}) \prod_{a=1}^{k-1} (l_{k-1,a} - l'_{ki}) \prod_{a=1, a\neq i}^{k} \frac{1}{l'_{ka} - l'_{ki}}.$$

We denote by $\Lambda \pm \delta_{ki}$ and $\Lambda + \delta'_{ki}$ the arrays obtained from Λ by replacing λ_{ki} and λ'_{ki} by $\lambda_{ki} \pm 1$ and $\lambda'_{ki} + 1$ respectively.

The following is our main theorem which will be proved in sections 2.4 and 2.5.

Theorem 2.1. *There exists a basis $\{\zeta_\Lambda\}$ of $V(\lambda)$ parametrized by the patterns Λ such that the action of the generators of $\mathfrak{g}_n$ is given by the formulae*

$$F_{kk}\zeta_\Lambda = \left(\sigma_k + 2\sum_{i=1}^{k} \lambda'_{ki} - \sum_{i=1}^{k} \lambda_{ki} - \sum_{i=1}^{k-1} \lambda_{k-1,i}\right)\zeta_\Lambda$$

$$F_{k-1,-k}\zeta_\Lambda = A_{k0}\zeta_\Lambda(k,0) + \sum_{i=1}^{k-1} A_{ki}\left(\frac{1}{l_{k-1,i} + \frac{1}{2}}\zeta^+_\Lambda(k,i) - \frac{1}{l_{k-1,i} - \frac{1}{2}}\zeta^-_\Lambda(k,i)\right).$$

Here the following notation has been used:

$$\zeta^-_\Lambda(k,i) = \zeta_{\Lambda - \delta_{k-1,i}}$$

$$\zeta^+_\Lambda(k,i) = \sum_{j=1}^{k} \sum_{m=1}^{k-1} B_{kj}(l_{k-1,i}) B_{k-1,m}(l_{k-1,i}) \zeta_{\Lambda + \delta'_{kj} + \delta_{k-1,i} + \delta'_{k-1,m}}$$

and $\zeta_\Lambda := 0$ if Λ is not a pattern. Furthermore,

$$\zeta_\Lambda(k,0) = (-1)^k \zeta_{\bar\Lambda} \qquad \textit{if} \quad \sigma_k = \sigma_{k-1} = 0$$

$$= \sum_{j=1}^{k} B_{kj}(l_{k-1,0}) \zeta_{\bar\Lambda + \delta'_{kj}} \qquad \textit{if} \quad \sigma_k = 1 \qquad \sigma_{k-1} = 0$$

$$= -\sum_{m=1}^{k-1} B_{k-1,m}(l_{k-1,0})\zeta_{\bar\Lambda+\delta'_{k-1,m}} \qquad \text{if} \quad \sigma_k = 0 \qquad \sigma_{k-1} = 1$$

$$= (-1)^{k-1} \sum_{j=1}^{k}\sum_{m=1}^{k-1} B_{kj}(l_{k-1,0})B_{k-1,m}(l_{k-1,0})\zeta_{\bar\Lambda+\delta'_{kj}+\delta'_{k-1,m}}$$
$$\text{if} \quad \sigma_k = \sigma_{k-1} = 1$$

where $\bar\Lambda$ is obtained from Λ by replacing σ_k and σ_{k-1} respectively with $\sigma_k + 1$ and $\sigma_{k-1} + 1$ (modulo 2). The action of $F_{k-1,k}$ is found from the relation

$$F_{k-1,k} = [\Phi_{k-1,-k}(u+2)\Phi_{-k,k} - \Phi_{-k,k}\Phi_{k-1,-k}(u)]_{u=0}$$

where the operator $\Phi_{-k,k}$ acts on the basis elements by the rule

$$\Phi_{-k,k}\zeta_\Lambda = \sum_{i=1}^{k} C_{ki}(F_{kk} - l'_{ki} + 1)\zeta_{\Lambda-\delta'_{ki}}$$

while the action of $\Phi_{k-1,-k}(u)$ is given by

$$\Phi_{k-1,-k}(u)\zeta_\Lambda = \frac{A_{k0}}{u + F_{kk} - \frac{3}{2}}\zeta_\Lambda(k,0)$$
$$+\sum_{i=1}^{k-1} A_{ki}\left(\frac{1}{(l_{k-1,i}+\frac{1}{2})(u + l_{k-1,i} + F_{kk} - 1)}\zeta^+_\Lambda(k,i)\right.$$
$$\left. - \frac{1}{(l_{k-1,i}-\frac{1}{2})(u - l_{k-1,i} + F_{kk} - 1)}\zeta^-_\Lambda(k,i)\right).$$

Remark. The image of ζ_Λ under the operator $\Phi_{k-1,-k}(u+2)\Phi_{-k,k} - \Phi_{-k,k}\Phi_{k-1,-k}(u)$ at $u = 0$ may be undefined for some patterns Λ. To get the action of $F_{k-1,k}$, one should first calculate its matrix elements in a 'generic' representation $V(\lambda)$ and then specialize the parameters; see section 2.5. For example, consider the case $n = 1$. It will be shown in section 2.4 that the basis vectors in $V(\lambda)$ are given by

$$\zeta_\Lambda = F_{10}^{\sigma_1}(F_{10}F_{0,-1})^{\lambda'_{11}-\lambda_{11}}\xi.$$

Furthermore, the operators $\Phi_{-1,1}$ and $\Phi_{0,-1}(u)$ are defined by

$$\Phi_{-1,1} = -\tfrac{1}{2}F_{01}^2$$

$$\Phi_{0,-1}(u) = F_{0,-1}\frac{1}{u + F_{11} - \frac{1}{2}}$$

see section 2.5. In the case $\lambda = (-\frac{1}{2})$ the basis of $V(\lambda)$ consists of two vectors ξ and $\xi' = F_{10}\xi$. Therefore, $\Phi_{-1,1}$ is the zero operator in $V(\lambda)$ while the image $\Phi_{0,-1}(0)\xi'$ is not defined. On the other hand, we find directly that $F_{01}\xi' = \frac{1}{2}\xi$.

2.2. *Transvector algebra* $\mathrm{Z}(\mathfrak{g}_n, \mathfrak{g}_{n-1})$

Consider the extension of the universal enveloping algebra $\mathrm{U}(\mathfrak{g}_n)$

$$\mathrm{U}'(\mathfrak{g}_n) = \mathrm{U}(\mathfrak{g}_n) \otimes_{\mathrm{U}(\mathfrak{h})} \mathrm{R}(\mathfrak{h})$$

where $\mathrm{R}(\mathfrak{h})$ is the field of fractions of the commutative algebra $\mathrm{U}(\mathfrak{h})$. Let J denote the left ideal in $\mathrm{U}'(\mathfrak{g}_n)$ generated by the elements F_{ij} with $-n < i < j < n$. The *transvector algebra* $\mathrm{Z}(\mathfrak{g}_n, \mathfrak{g}_{n-1})$ is the quotient algebra of the normalizer

$$\text{Norm J} = \{x \in \mathrm{U}'(\mathfrak{g}_n) | \mathrm{J}x \subseteq \mathrm{J}\}$$

modulo the two-sided ideal J [23]. It is an algebra over $\mathbb{C}$ and an $\mathrm{R}(\mathfrak{h})$-bimodule. Let $p = p_{n-1}$ denote the *extremal projection* for the Lie algebra $\mathfrak{g}_{n-1}$ [1, 23]. It satisfies the following (characteristic) relations:

$$F_{ij}\, p = p F_{ji} = 0 \qquad \text{for} \quad -n < i < j < n.$$

The projection p naturally acts in the space $\mathrm{U}'(\mathfrak{g}_n)/\mathrm{J}$ and its image coincides with $\mathrm{Z}(\mathfrak{g}_n, \mathfrak{g}_{n-1})$. The elements

$$pF_{ia} = -pF_{-a,-i} \qquad a = -n, n \qquad i = -n+1, \dots, n-1 \tag{2.4}$$

are generators of $\mathrm{Z}(\mathfrak{g}_n, \mathfrak{g}_{n-1})$ [23]. Set

$$f_0 = -\tfrac{1}{2} \qquad f_i = F_{ii} - i + \tfrac{1}{2} \qquad f_{-i} = -f_i$$

for $i = 1, \dots, n$; and set $f_i' = -f_{-i}$ for all i. The elements (2.4) can be given by the following explicit formulae (modulo J):

$$pF_{ia} = \sum_{i > i_1 > \cdots > i_s > -n} F_{ii_1} F_{i_1 i_2} \cdots F_{i_{s-1} i_s} F_{i_s a} \frac{1}{(f_i - f_{i_1}) \cdots (f_i - f_{i_s})}$$

where $s = 0, 1, \dots$. We shall use normalized generators of $\mathrm{Z}(\mathfrak{g}_n, \mathfrak{g}_{n-1})$ defined by

$$z_{ia} = pF_{ia}(f_i - f_{i-1}) \cdots (f_i - f_{-n+1}) \tag{2.5}$$

$$z_{ai} = pF_{ai}(f_i' - f_{i+1}') \cdots (f_i' - f_{n-1}'). \tag{2.6}$$

We obviously have $z_{ai} = (-1)^{n-i} z_{-i,-a}$. The elements z_{ia} satisfy certain quadratic relations [23]. We shall use the following ones below: for $a, b \in \{-n, n\}$ and $i + j \neq 0$ one has

$$z_{aj} z_{bi} (f_i' - f_j' + 1) = z_{bi} z_{aj} (f_i' - f_j') + z_{ai} z_{bj}. \tag{2.7}$$

In particular, z_{ai} and z_{aj} commute for $i + j \neq 0$. One easily verifies that z_{ai} and z_{bi} also commute for $i \neq 0$ and all a, b.

The elements z_{ia} and z_{ai} naturally act in the space $V(\lambda)^+$ and are called the *raising* and *lowering operators*. One has for $i = 1, \dots, n-1$:

$$z_{ia} : V(\lambda)^+_\mu \to V(\lambda)^+_{\mu+\delta_i} \qquad z_{ai} : V(\lambda)^+_\mu \to V(\lambda)^+_{\mu-\delta_i}$$

where $\mu \pm \delta_i$ is obtained from μ by replacing μ_i with $\mu_i \pm 1$. The operators z_{0a} preserve each subspace $V(\lambda)^+_\mu$.

We shall need the following element which can be checked to belong to the normalizer Norm J, and so it can be regarded as an element of the algebra $\mathrm{Z}(\mathfrak{g}_n, \mathfrak{g}_{n-1})$:

$$z_{n,-n} = \sum_{n > i_1 > \cdots > i_s > -n} F_{ni_1} F_{i_1 i_2} \cdots F_{i_s,-n} (f_n - f_{j_1}) \cdots (f_n - f_{j_k}) \tag{2.8}$$

where $s = 1, 2, \dots$ and $\{j_1, \dots, j_k\}$ is the complement to the subset $\{i_1, \dots, i_s\}$ in $\{-n+1, \dots, n-1\}$. The following relation is proved exactly as its C and D series counterparts [11, 12]: for $a = -n, n$

$$F_{n-1,a} = \sum_{i=-n+1}^{n-1} z_{n-1,i} z_{ia} \frac{1}{(f_i - f_{-n+1}) \cdots \wedge \cdots (f_i - f_{n-1})} \tag{2.9}$$

where $z_{n-1,n-1} := 1$ and the equalities are considered in $\mathrm{U}'(\mathfrak{g}_n)$ modulo the ideal J.

2.3. Yangians and twisted Yangians

Let us introduce the $\mathfrak{gl}(2)$-*Yangian* $\mathrm{Y}(2)$ and the (*orthogonal*) *twisted Yangian* $\mathrm{Y}^{+}(2)$; see [13] for more details. The Yangian $\mathrm{Y}(2)$ is the complex associative algebra with the generators $t_{ab}^{(1)}, t_{ab}^{(2)}, \ldots$ where $a, b \in \{-n, n\}$, and the defining relations

$$[t_{ab}(u), t_{cd}(v)] = \frac{1}{u-v}(t_{cb}(u)t_{ad}(v) - t_{cb}(v)t_{ad}(u)) \tag{2.10}$$

where

$$t_{ab}(u) := \delta_{ab} + t_{ab}^{(1)}u^{-1} + t_{ab}^{(2)}u^{-2} + \cdots \in \mathrm{Y}(2)[[u^{-1}]]. \tag{2.11}$$

Introduce the series $s_{ab}(u)$, $a, b \in \{-n, n\}$ by

$$s_{ab}(u) = t_{an}(u)t_{-b,-n}(-u) + t_{a,-n}(u)t_{-b,n}(-u).$$

Write $s_{ab}(u) = \delta_{ab} + s_{ab}^{(1)}u^{-1} + s_{ab}^{(2)}u^{-2} + \cdots$. The twisted Yangian $\mathrm{Y}^{+}(2)$ is defined as the subalgebra of $\mathrm{Y}(2)$ generated by the elements $s_{ab}^{(1)}, s_{ab}^{(2)}, \ldots$ where $a, b \in \{-n, n\}$. Also, $\mathrm{Y}^{+}(2)$ can be viewed as an abstract algebra with generators $s_{ab}^{(r)}$ and the following defining relations (see [13, section 3]):

$$\begin{aligned}[s_{ab}(u), s_{cd}(v)] &= \frac{1}{u-v}(s_{cb}(u)s_{ad}(v) - s_{cb}(v)s_{ad}(u)) \\ &\quad - \frac{1}{u+v}(s_{a,-c}(u)s_{-b,d}(v) - s_{c,-a}(v)s_{-d,b}(u)) \\ &\quad + \frac{1}{u^2-v^2}(s_{c,-a}(u)s_{-b,d}(v) - s_{c,-a}(v)s_{-b,d}(u))\end{aligned} \tag{2.12}$$

and

$$s_{-b,-a}(-u) = \frac{2u+1}{2u}s_{ab}(u) - \frac{1}{2u}s_{ab}(-u). \tag{2.13}$$

The Yangian $\mathrm{Y}(2)$ is a Hopf algebra with the coproduct

$$\Delta(t_{ab}(u)) = t_{an}(u) \otimes t_{nb}(u) + t_{a,-n}(u) \otimes t_{-n,b}(u). \tag{2.14}$$

The twisted Yangian $\mathrm{Y}^{+}(2)$ is a left coideal in $\mathrm{Y}(2)$ with

$$\Delta(s_{ab}(u)) = \sum_{c,d\in\{-n,n\}} t_{ac}(u)t_{-b,-d}(-u) \otimes s_{cd}(u). \tag{2.15}$$

Given a pair of complex numbers (α, β) such that $\alpha - \beta \in \mathbb{Z}_{+}$ we denote by $L(\alpha, \beta)$ the irreducible representation of the Lie algebra $\mathfrak{gl}(2)$ with the highest weight (α, β) with respect to the upper triangular Borel subalgebra; see the introduction. We have $\dim L(\alpha, \beta) = \alpha - \beta + 1$. We may regard $L(\alpha, \beta)$ as a $\mathrm{Y}(2)$-module by using the algebra homomorphism $\mathrm{Y}(2) \to \mathrm{U}(\mathfrak{gl}(2))$ given by

$$t_{ab}(u) \mapsto \delta_{ab} + E_{ab}u^{-1} \qquad a, b \in \{-n, n\}.$$

The coproduct (2.14) allows one to construct representations of $\mathrm{Y}(2)$ of the form

$$L = L(\alpha_1, \beta_1) \otimes \cdots \otimes L(\alpha_n, \beta_n).$$

Note that the generators $t_{ab}^{(r)}$ with $r > n$ act as zero operators in L. Therefore, the operators $T_{ab}(u) = u^n t_{ab}(u)$ are polynomials in u:

$$T_{ab}(u) = \delta_{ab}u^n + t_{ab}^{(1)}u^{n-1} + \cdots + t_{ab}^{(n)}.$$

For any $\gamma \in \mathbb{C}$ denote by $W(\gamma)$ the one-dimensional representation of $\mathrm{Y}^{+}(2)$ spanned by a vector w such that

$$s_{nn}(u)w = \frac{u+\gamma}{u+\frac{1}{2}}w \qquad s_{-n,-n}(u)w = \frac{u-\gamma+1}{u+\frac{1}{2}}w$$

and $s_{a,-a}(u)w = 0$ for $a = -n, n$. By (2.15) we can regard the tensor product $L \otimes W(\gamma)$ as a representation of $\mathrm{Y}^+(2)$. Representations of this type essentially exhaust all finite-dimensional irreducible representations of $\mathrm{Y}^+(2)$ [10]. The vector space isomorphism

$$L \otimes W(\gamma) \to L \qquad v \otimes w \mapsto v \qquad v \in L \tag{2.16}$$

provides $L \otimes W(\gamma)$ with an action of $\mathrm{Y}(2)$.

2.4. Construction of the basis

Introduce the following polynomials in u with coefficients in the transvector algebra $\mathrm{Z}(\mathfrak{g}_n, \mathfrak{g}_{n-1})$: for $a, b \in \{-n, n\}$

$$Z_{ab}(u) = -(\delta_{ab}(u - n + 1) + F_{ab}) \prod_{i=-n+1}^{n-1} (u + g_i) + \sum_{i=-n+1}^{n-1} z_{ai} z_{ib} \prod_{j=-n+1, j\neq i}^{n-1} \frac{u + g_j}{g_i - g_j} \tag{2.17}$$

where $g_i := f_i + \frac{1}{2}$ for all i.

Proposition 2.2. *The mapping*

$$s_{ab}(u) \mapsto -u^{-2n} Z_{ab}(u) \qquad a, b \in \{-n, n\} \tag{2.18}$$

defines an algebra homomorphism $\mathrm{Y}^+(2) \to \mathrm{Z}(\mathfrak{g}_n, \mathfrak{g}_{n-1})$.

Proof. One of the possible ways to prove the claim is to check directly that the relations (2.12) and (2.13) are satisfied with the $s_{ab}(u)$ respectively replaced with $Z_{ab}(u)$. Here one needs to use the quadratic relations in the transvector algebra $\mathrm{Z}(\mathfrak{g}_n, \mathfrak{g}_{n-1})$. In addition to (2.7) the relations which express $z_{ia} z_{bi}$ in terms of the $z_{bj} z_{ja}$ are needed; see [23].

Alternatively, we can follow the approach of [11, section 5] to construct first a homomorphism from $\mathrm{Y}^+(2)$ to the centralizer C_n of $\mathfrak{g}_{n-1}$ in $\mathrm{U}(\mathfrak{g}_n)$ and then calculate the images of the centralizer elements in the algebra $\mathrm{Z}(\mathfrak{g}_n, \mathfrak{g}_{n-1})$. The calculation is similar to that in the symplectic case [11]; see also [12]. We shall only give a few key formulae here. Introduce the $(2n + 1) \times (2n + 1)$ matrix $F = (F_{ij})$ whose ijth entry is the element $F_{ij} \in \mathfrak{g}_n$ and set

$$F(u) = 1 + \frac{F}{u + \frac{1}{2}}.$$

Denote by $\hat{F}(u)$ the corresponding *Sklyanin comatrix*; see [10, section 2]. The mapping

$$s_{ab}(u) \mapsto c(u)\hat{F}(-u + n - \tfrac{1}{2})_{ab} \qquad a, b \in \{-n, n\}, \tag{2.19}$$

where $c(u) = (1 - u^{-2})(1 - 4u^{-2}) \cdots (1 - (n-1)^2 u^{-2})$, defines an algebra homomorphism $\mathrm{Y}^+(2) \to \mathrm{C}_n$ [10, proposition 2.1]; cf [17]. Its composition with the natural homomorphism $\mathrm{C}_n \to \mathrm{Z}(\mathfrak{g}_n, \mathfrak{g}_{n-1})$ gives (2.18). □

As it follows from the branching rule (2.2), the space $V(\lambda)^+_\mu$ is nonzero only if there exists ν such that the inequalities (2.3) hold. We shall be assuming that this condition is satisfied. Proposition 2.2 allows one to equip $V(\lambda)^+_\mu$ with a structure of a $\mathrm{Y}^+(2)$-module defined via the homomorphism (2.18). The next theorem provides an identification of this module.

Theorem 2.3. *The* $\mathrm{Y}^+(2)$*-module* $V(\lambda)^+_\mu$ *is isomorphic to the direct sum of two irreducible submodules,* $V(\lambda)^+_\mu \simeq U \oplus U'$*, where*

$$U = L(0, \beta_1) \otimes L(\alpha_2, \beta_2) \otimes \cdots \otimes L(\alpha_n, \beta_n) \otimes W(\tfrac{1}{2}) \tag{2.20}$$

$$U' = L(-1, \beta_1) \otimes L(\alpha_2, \beta_2) \otimes \cdots \otimes L(\alpha_n, \beta_n) \otimes W(\tfrac{1}{2}) \tag{2.21}$$

if the λ_i are integers (it is supposed that $U' = \{0\}$ if $\beta_1 = 0$); or

$$U = L(-\tfrac{1}{2}, \beta_1) \otimes L(\alpha_2, \beta_2) \otimes \cdots \otimes L(\alpha_n, \beta_n) \otimes W(0) \tag{2.22}$$

$$U' = L(-\tfrac{1}{2}, \beta_1) \otimes L(\alpha_2, \beta_2) \otimes \cdots \otimes L(\alpha_n, \beta_n) \otimes W(1) \tag{2.23}$$

if the λ_i are half-integers, and the following notation is used:

$$\alpha_i = \min\{\lambda_{i-1}, \mu_{i-1}\} - i + 1 \qquad i = 2, \ldots, n$$

$$\beta_i = \max\{\lambda_i, \mu_i\} - i + 1 \qquad i = 1, \ldots, n$$

with $\mu_n := -\infty$. In particular, each of U and U' (and hence $V(\lambda)^+_\mu$) is equipped with an action of $\mathrm{Y}(2)$ defined by (2.16).

Proof. Consider the following two vectors in $V(\lambda)^+_\mu$:

$$\xi_\mu = \prod_{i=1}^{n-1} (z_{ni}^{\max\{\lambda_i,\mu_i\}-\mu_i} z_{i,-n}^{\max\{\lambda_i,\mu_i\}-\lambda_i})\xi \qquad \xi'_\mu = z_{n0}\xi_\mu. \tag{2.24}$$

Repeating the arguments of the proof of theorem 5.2 in [11] we can show that both ξ_μ and ξ'_μ are eigenvectors for $s_{nn}(u)$ and are annihilated by $s_{-n,n}(u)$. Namely,

$$s_{nn}(u)\xi_\mu = \mu(u)\xi_\mu \qquad s_{nn}(u)\xi'_\mu = (1 + u^{-1})\mu(u)\xi'_\mu \tag{2.25}$$

where

$$\mu(u) = (1 - \alpha_2 u^{-1}) \cdots (1 - \alpha_n u^{-1})(1 + \beta_1 u^{-1}) \cdots (1 + \beta_n u^{-1}).$$

This is proved simultaneously with the following relations by induction on the degree of the monomial in (2.24): for $i = 1, \ldots, n-1$

$$z_{in}\xi_\mu = -(m_i + \alpha_1) \cdots (m_i + \alpha_n)(m_i - \beta_1) \cdots (m_i - \beta_n)\xi_{\mu+\delta_i} \tag{2.26}$$

and

$$z_{-ni}\xi_\mu = -(m_i - \alpha_1 - 1) \cdots (m_i - \alpha_n - 1)(m_i + \beta_1 - 1) \cdots (m_i + \beta_n - 1)\xi_{\mu-\delta_i} \tag{2.27}$$

where $\alpha_1 = 0$ and $m_i = \mu_i - i + 1$ for $i = 1, \ldots, n-1$. Indeed, we note that if $\mu_i \geqslant \lambda_i$ then $z_{in}\xi_\mu = 0$ which is implied by (2.7). This agrees with (2.26) because in this case $\beta_i = m_i$. Now assume that $\mu_i < \lambda_i$. We have $z_{in}\xi_\mu = z_{in}z_{ni}\xi_{\mu+\delta_i}$ by (2.7). Formula (2.17) gives $z_{in}z_{ni} = z_{-n,-i}z_{-i,-n} = Z_{-n,-n}(-g_{-i})$. Further,

$$Z_{-n,-n}(-g_{-i})\xi_{\mu+\delta_i} = Z_{-n,-n}(m_i)\xi_{\mu+\delta_i}.$$

By proposition 2.2 and the symmetry relation (2.13) we can write

$$Z_{-n,-n}(m_i) = \frac{2m_i - 1}{2m_i} Z_{nn}(-m_i) + \frac{1}{2m_i} Z_{nn}(m_i).$$

By induction, $Z_{nn}(u)\xi_{\mu+\delta_i}$ can be found from (2.18) and (2.25) which gives (2.26). The proof of (2.27) is very similar. To prove (2.25) we apply the induction hypotheses to (2.26) and (2.27) and also use the relation

$$z_{0n}\xi_\mu = 0 \tag{2.28}$$

which is a consequence of (2.7). The relations

$$Z_{-n,n}(u)\xi_\mu = 0 \qquad Z_{-n,n}(u)\xi'_\mu = 0 \tag{2.29}$$

follow from (2.17), (2.26) and (2.27).

Both vectors ξ_μ and ξ'_μ are nonzero, except for the case $\beta_1 = 0$ where $\xi'_\mu = 0$. Indeed, applying appropriate operators z_{in} to ξ_μ or ξ'_μ repeatedly, we can obtain the highest vector ξ of $V(\lambda)$ with a nonzero coefficient. It follows from [10, corollary 6.6] that the tensor products (2.20)–(2.23) are irreducible representations of $Y^+(2)$. An easy calculation shows that the highest weights of the $Y^+(2)$-modules U and U' respectively coincide with the $Y^+(2)$-weights of the vectors ξ_μ and ξ'_μ. So, U and U' are respectively isomorphic to quotients of the $Y^+(2)$-submodules in $V(\lambda)^+_\mu$ generated by ξ_μ and ξ'_μ. On the other hand, the branching rule (2.2) implies that

$$\dim V(\lambda)^+_\mu = \dim U + \dim U'.$$

Therefore, to complete the proof of the theorem we need to show that the $Y^+(2)$-submodules generated by ξ_μ and ξ'_μ are disjoint. For this we employ a contravariant bilinear form $\langle , \rangle$ on $V(\lambda)$ uniquely determined by the conditions

$$\langle \xi, \xi \rangle = 1 \qquad \langle F_{ij}\eta, \zeta \rangle = \langle \eta, F_{ji}\zeta \rangle \qquad \eta, \zeta \in V(\lambda).$$

One easily shows that its restriction to the subspace $V(\lambda)^+_\mu$ is nondegenerate. Therefore, our claim will follow from the fact that the submodules generated by ξ_μ and ξ'_μ are orthogonal to each other with respect to $\langle , \rangle$. Given an operator A in $V(\lambda)^+$ we denote by A^* its adjoint operator with respect to the form:

$$\langle A\eta, \zeta \rangle = \langle \eta, A^*\zeta \rangle.$$

Since the extremal projection p is stable with respect to the anti-involution $F_{ij} \mapsto F_{ji}$ [23] we derive that $(pF_{ia})^* = pF_{ai}$ for $a = -n, n$ and $i = -n+1, \ldots, n-1$. Therefore, $z^*_{ia} = z_{ai} \cdot c$, where c is an element of $\mathrm{R}(\mathfrak{h}_{n-1})$ which can be found from (2.5) and (2.6). This also implies that $(z_{ai}z_{ib})^* = z_{bi}z_{ia}$ and hence

$$Z_{ab}(u)^* = Z_{ba}(u) \tag{2.30}$$

see (2.17). By proposition 2.2 and the Poincaré–Birkhoff–Witt theorem for the twisted Yangians [13, remark 3.14], every element of the $Y^+(2)$-submodules generated by ξ_μ and ξ'_μ can be written as a linear combination of vectors of the following form, respectively:

$$Z_{n,-n}(u_1)\cdots Z_{n,-n}(u_k)\xi_\mu \qquad \text{or} \qquad Z_{n,-n}(v_1)\cdots Z_{n,-n}(v_l)\xi'_\mu$$

where the u_i and v_i are complex parameters. Therefore, by (2.29) and (2.30) to prove that the submodules are orthogonal it now suffices to show that $\langle \xi_\mu, \xi'_\mu \rangle = 0$. But this follows from (2.28). □

Remark. Using Weyl's formula for the dimension of $V(\lambda)$ one can slightly modify the proof of theorem 2.3 so that the branching rule (2.2) would not be used but follows from the theorem; cf [11, 12].

It follows from (2.13) that the series $s_{n,-n}(u)$ is even in u, and so is the polynomial $Z_{n,-n}(u)$; see proposition 2.2. On the other hand, (2.17) implies that $Z_{n,-n}(-g_i) = z_{ni}z_{i,-n}$ for $i = 1, \ldots, n-1$. Moreover, $Z_{n,-n}(-g_n) = z_{n,-n}$ which follows from (2.8). Since $Z_{n,-n}(u)$ is a polynomial in u^2 of degree $n-1$, by the Lagrange interpolation formula $Z_{n,-n}(u)$ can also be given by

$$Z_{n,-n}(u) = \sum_{i=1}^{n} z_{ni}z_{i,-n} \prod_{j=1, j\neq i}^{n} \frac{u^2 - g_j^2}{g_i^2 - g_j^2}. \tag{2.31}$$

Remark. To make the above evaluation $Z_{n,-n}(-g_i)$ well defined we agree to consider the series $Z_{ab}(u)$ with $a, b \in \{-n, n\}$ as elements of the *right* module over the field of rational functions in $g_1, \ldots, g_n, u$ generated by monomials in the z_{ia}.

Given ν such that the conditions (2.3) are satisfied, set

$$\gamma_i = \nu_i - i + 1 \qquad l_i = \lambda_i - i + 1 \qquad i \geqslant 1$$

and introduce the vectors

$$\xi_{\nu\mu} = \begin{cases} \prod_{i=1}^{n} Z_{n,-n}(\gamma_i - 1) \cdots Z_{n,-n}(\beta_i + 1) Z_{n,-n}(\beta_i)\xi_\mu & \text{if} \quad \sigma = 0 \\ \prod_{i=1}^{n} Z_{n,-n}(\gamma_i - 1) \cdots Z_{n,-n}(\beta_i + 1) Z_{n,-n}(\beta_i)\xi'_\mu & \text{if} \quad \sigma = 1. \end{cases}$$

Using (2.31) and (2.7) we get an equivalent expression; cf [11, section 6]:

$$\xi_{\nu\mu} = z_{n0}^{\sigma} \prod_{i=1}^{n-1} z_{ni}^{\nu_i - \mu_i} z_{i,-n}^{\nu_i - \lambda_i} \cdot \prod_{k=l_n}^{\gamma_n - 1} Z_{n,-n}(k)\xi.$$

Proposition 2.4. *The vectors $\xi_{\nu\mu}$ with ν satisfying (2.3) form a basis of $V(\lambda)^+_\mu$.*

Proof. Due to theorem 2.3 it suffices to show that the vectors $\xi_{\nu\mu}$ with $\sigma = 0$ form a basis of the subspace U of $V(\lambda)^+_\mu$ while those with $\sigma = 1$ form a basis in U'. Let us write each of the tensor products in (2.20)–(2.23) in the form

$$L(\alpha_1, \beta_1) \otimes L(\alpha_2, \beta_2) \otimes \cdots \otimes L(\alpha_n, \beta_n) \otimes W(-\alpha_0). \tag{2.32}$$

Regarding this as a Y(2)-module defined by (2.16) we can construct a Gelfand–Tsetlin-type basis in this module as follows. Set

$$\tilde{\zeta}_{\nu\mu} = \prod_{i=1}^{n} T_{n,-n}(-\gamma_i + 1) \cdots T_{n,-n}(-\beta_i - 1) T_{n,-n}(-\beta_i) z_{n0}^{\sigma}\xi_\mu.$$

The vectors $\tilde{\zeta}_{\nu\mu}$ with ν satisfying (2.3) form a basis in the Y(2)-module (2.32); see [9, 16, 19]. Furthermore, we have

$$\begin{aligned} T_{nn}(u)\tilde{\zeta}_{\nu\mu} &= (u + \gamma_1) \cdots (u + \gamma_n)\tilde{\zeta}_{\nu\mu} \\ T_{n,-n}(-\gamma_i)\tilde{\zeta}_{\nu\mu} &= \tilde{\zeta}_{\nu+\delta_i,\mu}. \end{aligned} \tag{2.33}$$

We have the following equality of operators in the space (2.32):

$$Z_{n,-n}(u) = \frac{(u - \alpha_0)T_{n,-n}(-u)T_{nn}(u) + (u + \alpha_0)T_{n,-n}(u)T_{nn}(-u)}{(-1)^{n+1}u}$$

which is easily derived from (2.10), (2.15) and (2.18). Therefore, by (2.33)

$$Z_{n,-n}(\gamma_i)\tilde{\zeta}_{\nu\mu} = -2(\alpha_0 - \gamma_i) \prod_{a=1, a\neq i}^{n} (-\gamma_a - \gamma_i)\tilde{\zeta}_{\nu+\delta_i,\mu}.$$

This shows that for each ν the vectors $\xi_{\nu\mu}$ and $\tilde{\zeta}_{\nu\mu}$ coincide up to a nonzero factor. □

We shall use the following normalized basis vectors:

$$\zeta_{\nu\mu} = \prod_{1\leqslant i<j\leqslant n} (-\gamma_i - \gamma_j)!\ \xi_{\nu\mu}.$$

The following formulae for the action of the generators of the Yangian Y(2) in the basis $\{\zeta_{\nu\mu}\}$ follow from the above proof: for $i = 1, \ldots, n$

$$\begin{aligned} T_{nn}(u)\zeta_{\nu\mu} &= (u + \gamma_1) \cdots (u + \gamma_n)\zeta_{\nu\mu} \\ T_{n,-n}(-\gamma_i)\zeta_{\nu\mu} &= \frac{1}{2(\gamma_i - \alpha_0)}\zeta_{\nu+\delta_i,\mu} \\ T_{-n,n}(-\gamma_i)\zeta_{\nu\mu} &= 2\prod_{k=0}^{n}(\alpha_k - \gamma_i + 1)\prod_{k=1}^{n}(\beta_k - \gamma_i)\zeta_{\nu-\delta_i,\mu} \end{aligned} \tag{2.34}$$

cf [11] and [12].

Given a pattern Λ (see section 2.1) we introduce the vector

$$\xi_\Lambda = \prod_{k=1,\ldots,n}^{\rightarrow} \left(z_{k0}^{\sigma_k} \cdot \prod_{i=1}^{k-1} z_{ki}^{\lambda'_{ki}-\lambda_{k-1,i}} z_{i,-k}^{\lambda'_{ki}-\lambda_{ki}} \cdot \prod_{q=l_{kk}}^{l'_{kk}-1} Z_{k,-k}(q+\tfrac{1}{2}) \right) \xi$$

and set

$$\zeta_\Lambda = N_\Lambda \xi_\Lambda \qquad N_\Lambda = \prod_{k=2}^{n} \prod_{1\leqslant i<j\leqslant k} (-l'_{ki} - l'_{kj} - 1)!.$$

The following proposition is implied by the branching rule (2.2) and proposition 2.4.

Proposition 2.5. *The vectors ζ_Λ parametrized by the patterns Λ form a basis of the representation $V(\lambda)$.*

2.5. Matrix element formulae

Introduce the following elements of $\mathrm{U}(\mathfrak{g}_n)$:

$$\Phi_{-k,k} = \sum_{i=1}^{k-1} F_{-k,i} F_{ik} - \tfrac{1}{2} F_{0k}^2 \qquad k = 1, \ldots, n.$$

We shall find the action of $\Phi_{-k,k}$ in the basis $\{\zeta_\Lambda\}$, which will be used later on. Since $\Phi_{-k,k}$ commutes with the subalgebra $\mathfrak{g}_{k-1}$ it suffices to consider the case $k = n$. The image of $2\Phi_{-n,n}$ under the natural homomorphism $\pi : \mathrm{C}_n \to \mathrm{Z}(\mathfrak{g}_n, \mathfrak{g}_{n-1})$ coincides with the coefficient at u^{2n-2} of the polynomial $Z_{-n,n}(u)$; see the proof of proposition 2.2. The following equality of operators in (2.32) is obtained from (2.10), (2.15) and (2.18):

$$Z_{-n,n}(u) = \frac{(u-\alpha_0)T_{-n,-n}(-u)T_{-n,n}(u) + (u+\alpha_0)T_{-n,-n}(u)T_{-n,n}(-u)}{(-1)^{n+1}u}.$$

Therefore,

$$\Phi_{-n,n} = -t_{-n,n}^{(2)} + t_{-n,n}^{(1)} t_{-n,-n}^{(1)} + (1+\alpha_0) t_{-n,n}^{(1)}. \tag{2.35}$$

The image of $s_{nn}^{(1)}$ under the homomorphism (2.19) is F_{nn}. On the other hand, by (2.15) we have

$$s_{nn}^{(1)} = t_{nn}^{(1)} - t_{-n,-n}^{(1)} - \alpha_0 - \tfrac{1}{2}$$

as operators in the space (2.32). Therefore, (2.35) can be written as

$$\Phi_{-n,n} = -t_{-n,n}^{(2)} + t_{-n,n}^{(1)} t_{nn}^{(1)} - (F_{nn} + \tfrac{3}{2}) t_{-n,n}^{(1)}.$$

Finally, relations (2.34) imply that

$$\Phi_{-n,n} \zeta_{\nu\mu} = \sum_{i=1}^{n} \theta_i (F_{nn} - \gamma_i + \tfrac{3}{2}) \zeta_{\nu-\delta_i,\mu} \tag{2.36}$$

where

$$\theta_i = -2 \prod_{k=0}^{n} (\alpha_k - \gamma_i + 1) \prod_{k=1}^{n} (\beta_k - \gamma_i) \prod_{j=1, j\neq i}^{n} (\gamma_j - \gamma_i)^{-1}.$$

Using the notation in (2.32) we can also write this as

$$\theta_i = (2\gamma_i - 1)(1 - \sigma - \gamma_i) \prod_{k=1}^{n} (l_k - \gamma_i) \prod_{k=1}^{n-1} (m_k - \gamma_i) \prod_{j=1, j\neq i}^{n} (\gamma_j - \gamma_i)^{-1}.$$

The action of F_{nn} in $V(\lambda)^+_\mu$ is immediately found so that

$$F_{nn}\xi_{\nu\mu} = \left(\sigma + 2\sum_{i=1}^{n}\nu_i - \sum_{i=1}^{n}\lambda_i - \sum_{i=1}^{n-1}\mu_i\right)\xi_{\nu\mu}.$$

The operator $F_{n-1,-n}$ preserves the subspace of $\mathfrak{g}_{n-2}$-highest vectors in $V(\lambda)$. Therefore, it suffices to calculate its action on the basis vectors of the form

$$\xi_{\nu\mu\nu'} = X_{\mu\nu'}\xi_{\nu\mu}$$

where $X_{\mu\nu'}$ denotes the operator

$$X_{\mu\nu'} = z_{n-1,0}^{\sigma'}\prod_{i=1}^{n-2} z_{n-1,i}^{\nu_i'-\mu_i'} z_{i,-n+1}^{\nu_i'-\mu_i} \cdot \prod_{a=m_{n-1}}^{\gamma_{n-1}'-1} Z_{n-1,-n+1}(a).$$

Here we assume that the conditions (2.3) are satisfied with λ, ν, μ respectively replaced by μ, ν', μ'; we have used the notation $\gamma_i' = \nu_i' - i + 1$. The operator $F_{n-1,-n}$ is permutable with the elements $z_{n-1,i}$, $z_{i,-n+1}$ and $Z_{n-1,-n+1}(u)$ which follows from their explicit formulae. Hence, we can write

$$F_{n-1,-n}\xi_{\nu\mu\nu'} = X_{\mu\nu'}F_{n-1,-n}\xi_{\nu\mu}.$$

Let us apply (2.9) with $a = -n$. We have

$$f_i\xi_{\nu\mu} = (m_i - \tfrac{1}{2})\xi_{\nu\mu} \qquad f_{-i}\xi_{\nu\mu} = (-m_i + \tfrac{1}{2})\xi_{\nu\mu}$$

for $i \geqslant 1$. Recall also that $f_0 = -\frac{1}{2}$. We now need to express

$$X_{\mu\nu'}z_{n-1,i}z_{i,-n}\xi_{\nu\mu} \qquad i = -n+1, \ldots, n-1$$

as a linear combination of the vectors $\xi_{\nu\mu\nu'}$. Suppose first that $i \geqslant 1$. Assuming that $\nu_i - \mu_i \geqslant 1$ we obtain from (2.7)

$$z_{i,-n}\xi_{\nu\mu} = z_{i,-n}z_{ni}\xi_{\nu,\mu+\delta_i}. \tag{2.37}$$

By (2.17) this equals

$$z_{i,-n}z_{ni}\xi_{\nu,\mu+\delta_i} = z_{n,-i}z_{-i,-n}\xi_{\nu,\mu+\delta_i} = Z_{n,-n}(-g_{-i})\xi_{\nu,\mu+\delta_i}.$$

We have $-g_{-i}\xi_{\nu,\mu+\delta_i} = m_i\xi_{\nu,\mu+\delta_i}$. Since $Z_{n,-n}(\gamma_p)\xi_{\nu,\mu+\delta_i} = \xi_{\nu+\delta_p,\mu+\delta_i}$ for each $p = 1, \ldots, n$, we obtain from the Lagrange interpolation formula that (2.37) takes the form

$$Z_{n,-n}(m_i)\xi_{\nu,\mu+\delta_i} = \sum_{p=1}^{n}\prod_{a=1,a\neq p}^{n}\frac{m_i^2 - \gamma_a^2}{\gamma_p^2 - \gamma_a^2}\xi_{\nu+\delta_p,\mu+\delta_i}.$$

Furthermore, for $i \geqslant 1$

$$z_{-i,-n}\xi_{\nu\mu} = (-1)^{n-i}z_{ni}\xi_{\nu\mu} = (-1)^{n-i}\xi_{\nu,\mu-\delta_i}.$$

Consider now the vector $z_{0,-n}\xi_{\nu\mu}$. If $\sigma = 0$ then it equals

$$z_{0,-n}\xi_{\nu\mu} = (-1)^n z_{n0}\xi_{\nu\mu} = (-1)^n\xi_{\overline{\nu}\mu}$$

where $\overline{\nu} = (\sigma + 1, \nu_1, \ldots, \nu_n)$ (addition modulo 2). If $\sigma = 1$ then

$$z_{0,-n}\xi_{\nu\mu} = z_{n0}z_{0,-n}\xi_{\overline{\nu}\mu}$$

which coincides with $Z_{n,-n}(-g_0)\xi_{\overline{\nu}\mu}$, where $g_0 = 0$. Using again the Lagrange interpolation formula we find that this equals

$$Z_{n,-n}(m_0)\xi_{\overline{\nu}\mu} = \sum_{p=1}^{n}\prod_{a=1,a\neq p}^{n}\frac{m_0^2 - \gamma_a^2}{\gamma_p^2 - \gamma_a^2}\xi_{\overline{\nu}+\delta_p,\mu}$$

with $m_0 = 0$. The operator $X_{\mu\nu'} z_{n-1,i}$ is transformed in exactly the same manner; cf [11]. Combining the results we obtain

$$F_{n-1,-n}\xi_{\nu\mu\nu'} = A_0\xi(0) + \sum_{i=1}^{n-1} A_i\left(\frac{1}{m_i}\xi^+(i) - \frac{1}{m_i-1}\xi^-(i)\right)$$

where

$$A_i = \prod_{a=1,a\neq i}^{n-1}\frac{1}{m_i-m_a}\prod_{a=1}^{n-1}\frac{1}{m_i+m_a-1}.$$

Furthermore,

$$\xi^-(i) = \xi_{\nu,\mu-\delta_i,\nu'}$$

$$\xi^+(i) = \sum_{p=1}^{n}\sum_{q=1}^{n-1}\prod_{a=1,a\neq p}^{n}\frac{m_i^2-\gamma_a^2}{\gamma_p^2-\gamma_a^2}\prod_{a=1,a\neq q}^{n-1}\frac{m_i^2-\gamma_a'^2}{\gamma_q'^2-\gamma_a'^2}\xi_{\nu+\delta_p,\mu+\delta_i,\nu'+\delta_q}$$

and

$$\begin{aligned}
\xi(0) &= (-1)^n\xi_{\overline{\nu}\mu\overline{\nu}'} \qquad \text{if} \quad \sigma=\sigma'=0\\
&= \sum_{p=1}^{n}\prod_{a=1,a\neq p}^{n}\frac{m_0^2-\gamma_a^2}{\gamma_p^2-\gamma_a^2}\xi_{\overline{\nu}+\delta_p,\mu,\overline{\nu}'} \qquad \text{if} \quad \sigma=1 \qquad \sigma'=0\\
&= -\sum_{q=1}^{n-1}\prod_{a=1,a\neq q}^{n-1}\frac{m_0^2-\gamma_a'^2}{\gamma_q'^2-\gamma_a'^2}\xi_{\overline{\nu},\mu,\overline{\nu}'+\delta_q} \qquad \text{if} \quad \sigma=0 \qquad \sigma'=1\\
&= (-1)^{n-1}\sum_{p=1}^{n}\sum_{q=1}^{n-1}\prod_{a=1,a\neq p}^{n}\frac{m_0^2-\gamma_a^2}{\gamma_p^2-\gamma_a^2}\prod_{a=1,a\neq q}^{n-1}\frac{m_0^2-\gamma_a'^2}{\gamma_q'^2-\gamma_a'^2}\xi_{\overline{\nu}+\delta_p,\mu,\overline{\nu}'+\delta_q}\\
&\quad \text{if} \quad \sigma=\sigma'=1
\end{aligned}$$

with $\overline{\nu}' = (\sigma'+1, \nu_1', \ldots, \nu_{n-1}')$ (addition modulo 2). We now compute the action of $F_{n-1,n}$. In the formula (2.9) with $a = n$, replace the operators z_{in} with the following expression: for $i = -n+1, \ldots, n-1$

$$z_{in} = [z_{i,-n}, \Phi_{-n,n}]\frac{1}{f_i + F_{nn}} \tag{2.38}$$

and then use the formulae for the action of $z_{i,-n}$ and $\Phi_{-n,n}$; see (2.36). More precisely, we regard (2.38) as a relation in the transvector algebra $\mathrm{Z}(\mathfrak{g}_n, \mathfrak{g}_{n-1})$ which can be proved as follows. First, we calculate the commutator $[F_{i,-n}, \Phi_{-n,n}]$ in $\mathrm{U}(\mathfrak{g}_n)$ then consider it modulo the ideal J and apply the extremal projection p (see section 2).

We have $\Phi_{-n,n}F_{nn} = (F_{nn}+2)\Phi_{-n,n}$, and so (2.9) and (2.38) imply that

$$F_{n-1,n}\xi_{\nu\mu\nu'} = X_{\mu\nu'}\left(\Phi_{n-1,-n}(2)\Phi_{-n,n} - \Phi_{-n,n}\Phi_{n-1,-n}(0)\right)\xi_{\nu\mu} \tag{2.39}$$

where

$$\Phi_{n-1,-n}(u) = \sum_{i=-n+1}^{n-1} z_{n-1,i}z_{i,-n}\prod_{a=-n+1,a\neq i}^{n-1}\frac{1}{f_i-f_a}\cdot\frac{1}{u+f_i+F_{nn}}. \tag{2.40}$$

The action of $\Phi_{n-1,-n}(u)$ is found exactly as that of $F_{n-1,-n}$. Formula (2.39) is valid provided the denominators in (2.40) do not vanish. However, since (2.38) holds in the transvector algebra $\mathrm{Z}(\mathfrak{g}_n, \mathfrak{g}_{n-1})$, the relation (2.39) holds for generic parameters ν, μ and ν' which allows one to get explicit formulae for all matrix elements of $F_{n-1,n}$.

Finally, the proof of theorem 2.1 is completed by rewriting the above formulae for the action of the generators in terms of the parameters σ_k, l_{ki} and l'_{ki} of the patterns Λ. The parameters l_i, γ_i, m_i are replaced by

$$l_i \mapsto l_{ki} + \tfrac{1}{2} \qquad \gamma_i \mapsto l'_{ki} + \tfrac{1}{2} \qquad m_i \mapsto l_{k-1,i} + \tfrac{1}{2}.$$

Remark. There is another way of calculating the matrix elements of $F_{k-1,k}$ based on the formulae for the action of $T_{-n,-n}(u)$ in the basis $\{\zeta_{\nu\mu}\}$; see [12].

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